

On Fromm-Hill integral

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University of Windsor, PSAS 2008

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Motivation

For few-body systems,

$$H = H_0 + H_{rel} + H_{QED} + \dots$$

with H_0 a non-relativistic Hamiltonian and some corrections.

From Rayleigh-Ritz variational principle, we have a generalized eigenvalue problem

$$\underline{H_0} \underline{c} = \lambda \underline{O} \underline{c},$$

$$H_{0ij} = \langle \psi_i | H_0 | \psi_j \rangle,$$

$$O_{ij} = \langle \psi_i | \psi_j \rangle,$$

ie, we solve the eigenvalue problem for H_0 first.

Hylleraas-type wave functions are always chosen to be trial basis.

The basic integral

$$\begin{aligned} & J(n_1, n_2, n_3, n_{12}, n_{23}, n_{31}; \alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) \\ &= \int r_1^{n_1-1} r_2^{n_2-1} r_3^{n_3-1} r_{12}^{n_{12}-1} r_{23}^{n_{23}-1} r_{31}^{n_{31}-1} \\ & \times \exp(-\alpha_1 r_1 - \alpha_2 r_2 - \alpha_3 r_3 - \alpha_{12} r_{12} - \alpha_{23} r_{23} - \alpha_{31} r_{31}) \\ & \times d^3 \mathbf{r}_1 d^3 \mathbf{r}_2 d^3 \mathbf{r}_3. \end{aligned} \tag{1}$$

The generating integral

$$\begin{aligned} & I(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) \\ &= \int \frac{\exp(-\alpha_1 r_1 - \alpha_2 r_2 - \alpha_3 r_3 - \alpha_{12} r_{12} - \alpha_{23} r_{23} - \alpha_{31} r_{31})}{r_1 r_2 r_3 r_{12} r_{23} r_{31}} \\ & \times d^3 \mathbf{r}_1 d^3 \mathbf{r}_2 d^3 \mathbf{r}_3. \end{aligned} \tag{2}$$

"Generating" means that (1) can be derived from (2) by differentiation:

$$\begin{aligned} & J(n_1, n_2, n_3, n_{12}, n_{23}, n_{31}; \alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) \\ &= \left(-\frac{\partial}{\partial \alpha_1}\right)^{n_1} \left(-\frac{\partial}{\partial \alpha_2}\right)^{n_2} \left(-\frac{\partial}{\partial \alpha_3}\right)^{n_3} \left(-\frac{\partial}{\partial \alpha_{12}}\right)^{n_{12}} \\ & \quad \times \left(-\frac{\partial}{\partial \alpha_{23}}\right)^{n_{23}} \left(-\frac{\partial}{\partial \alpha_{31}}\right)^{n_{31}} I(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) \end{aligned} \quad (3)$$

- Nonrelativistic case: the integers $n_1, n_2, n_3, n_{12}, n_{23}$ and n_{31} are all nonnegative.
- Leading relativistic corrections: some of the integers are equal to -1.
- QED corrections, etc.

Remiddi-Pachucki-Puchalski method

Setting $\alpha_{\mu\nu} = 0$, the basic integral (1) is simplified to

$$\begin{aligned} & \tilde{J}(n_1, n_2, n_3, n_{12}, n_{23}, n_{31}; \alpha_1, \alpha_2, \alpha_3) \\ &= \int r_1^{n_1-1} r_2^{n_2-1} r_3^{n_3-1} r_{12}^{n_{12}-1} r_{23}^{n_{23}-1} r_{31}^{n_{31}-1} \\ & \quad \times \exp(-\alpha_1 r_1 - \alpha_2 r_2 - \alpha_3 r_3) d^3 \mathbf{r}_1 d^3 \mathbf{r}_2 d^3 \mathbf{r}_3. \end{aligned} \quad (4)$$

Consider the following "generating" integral,

$$\begin{aligned} & \tilde{I}(\alpha_1, \alpha_2, \alpha_3) \\ &= \int \frac{\exp(-\alpha_1 r_1 - \alpha_2 r_2 - \alpha_3 r_3)}{r_1 r_2 r_3 r_{12} r_{23} r_{31}} d^3 \mathbf{r}_1 d^3 \mathbf{r}_2 d^3 \mathbf{r}_3. \end{aligned} \quad (5)$$

E. Remiddi (PRA, **44**,9(1991)) derived a closed-form formula of the above integral

$$\begin{aligned}
& \tilde{I}(\alpha_1, \alpha_2, \alpha_3) \\
&= \frac{32\pi^3}{\alpha_1\alpha_2\alpha_3} \left\{ \ln \left[\frac{\alpha_1 + \alpha_2}{\alpha_3} \right] \ln \left[\frac{\alpha_1 + \alpha_2 + \alpha_3}{\alpha_1 + \alpha_2} \right] - \text{Li}_2 \left[-\frac{\alpha_3}{\alpha_1 + \alpha_2} \right] \right. \\
&\quad - \text{Li}_2 \left[1 - \frac{\alpha_3}{\alpha_1 + \alpha_2} \right] + \ln \left[\frac{\alpha_3 + \alpha_1}{\alpha_2} \right] \ln \left[\frac{\alpha_1 + \alpha_2 + \alpha_3}{\alpha_3 + \alpha_1} \right] \\
&\quad - \text{Li}_2 \left[-\frac{\alpha_2}{\alpha_3 + \alpha_1} \right] - \text{Li}_2 \left[1 - \frac{\alpha_2}{\alpha_3 + \alpha_1} \right] \\
&\quad + \ln \left[\frac{\alpha_2 + \alpha_3}{\alpha_1} \right] \ln \left[\frac{\alpha_1 + \alpha_2 + \alpha_3}{\alpha_2 + \alpha_3} \right] \\
&\quad \left. - \text{Li}_2 \left[-\frac{\alpha_1}{\alpha_2 + \alpha_3} \right] - \text{Li}_2 \left[1 - \frac{\alpha_1}{\alpha_2 + \alpha_3} \right] \right\}.
\end{aligned}$$

The closed-form result for the integral (4) with $n_1 = 1$, $n_2 = 1$, $n_3 = 1$, $n_{12} = 2$, $n_{23} = 0$ and $n_{31} = 2$ was also derived in Remiddi's paper.

With the closed-form expression for $\tilde{I}(\alpha_1, \alpha_2, \alpha_3)$, one can easily derive that

$$\tilde{J}(n_1, n_2, n_3, 0, 0, 0) = \left(-\frac{\partial}{\partial \alpha_1}\right)^{n_1} \left(-\frac{\partial}{\partial \alpha_2}\right)^{n_2} \left(-\frac{\partial}{\partial \alpha_3}\right)^{n_3} \times \tilde{I}(\alpha_1, \alpha_2, \alpha_3). \quad (6)$$

Now a series of questions arise:

- Can it produce the expression for the general integral $\tilde{J}(n_1, n_2, n_3, n_{12}, n_{23}, n_{31})$ with n_{12}, n_{23}, n_{31} being arbitrary nonnegative integers?
- And how?

$$\begin{aligned}
& \tilde{J}(n_1, n_2, n_3 + 1, n_{12}, n_{23}, n_{31}) \\
&= \frac{1}{\alpha_1 \alpha_2 \alpha_3} \left\{ (n_{23} - 1) n_{23} n_1 \tilde{J}(n_1 - 1, n_2 + 1, n_3, n_{12}, n_{23} - 2, n_{31}) \right. \\
&+ (n_{31} - 1) n_{31} n_2 \tilde{J}(n_1 + 1, n_2 - 1, n_3, n_{12}, n_{23}, n_{31} - 2) \\
&\left. - (n_{12} - 1) n_{12} n_1 \tilde{J}(n_1 - 1, n_2 - 1, n_3, n_{12}, n_{23}, n_{31}) + \dots \right\}. \tag{7}
\end{aligned}$$

Two other recursions for $\tilde{J}(n_1 + 1, n_2, n_3, n_{12}, n_{23}, n_{31})$ and $\tilde{J}(n_1, n_2 + 1, n_3, n_{12}, n_{23}, n_{31})$ can be obtained by symmetry.

$$\begin{aligned}
& \tilde{J}(0, 0, 0, n_{12} + 2, n_{23}, n_{31}) \\
&= \frac{1 + n_{12}}{2} \left\{ \frac{1}{\alpha_1^2} \left[\frac{(n_{31} - 1)n_{31}}{1 + n_{12}} \tilde{J}(0, 0, 0, n_{12} + 2, n_{23}, n_{31} - 2) \right. \right. \\
&+ (n_{12} + n_{23} + 2n_{31}) \tilde{J}(0, 0, 0, n_{12}, n_{23}, n_{31}) \\
&+ \left. \left. \frac{(n_{31} - 1)n_{31}}{1 + n_{23}} \tilde{J}(0, 0, 0, n_{12} + 2, n_{23}, n_{31} - 2) + \dots \right] \right. \\
&+ \left. \frac{1}{\alpha_2^2} [\dots] - \frac{\alpha_3^2}{\alpha_1^2 \alpha_2^2} [\dots] \right\}.
\end{aligned} \tag{8}$$

Two other recursions for $\tilde{J}(0, 0, 0, n_{12}, n_{23} + 2, n_{31})$ and $\tilde{J}(0, 0, 0, n_{12}, n_{23}, n_{31} + 2)$ can be obtained by symmetry.

Using the recursions given above, the boundary terms and some other conditions, one can evaluate the integral $\tilde{J}(n_1, n_2, n_3, n_{12}, n_{23}, n_{31})$ analytically.

(M. Puchalski, K. Pachucki and E. Remiddi, PRA **70**, 032502 (2004))

Furthermore, by introducing and developing some numerical methods, some evaluations for [relativistic corrections](#) are obtained.

(K. Pachucki et al, PRA **71**, 032514 (2005); K. Puchalski et al, PRA **73**, 022503 (2006))

Fromm-Hill integral

An analytic expression for the generating integral (2) is given by (D. M. Fromm and R. N. Hill, PRA **36**, 3(1987))

$$I(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) = \frac{16\pi^3}{\sigma} \left[\sum_{j=1}^3 u(\beta_0^{(0)} \beta_0^{(j)}) + \sum_{j=0}^3 \sum_{k=0}^3 v(\gamma_k^{(j)} / \sigma) \right], \quad (9)$$

where $u(z) = \text{Li}_2(z) - \text{Li}_2(1/z)$ and

$$v(z) = \frac{1}{2} \text{Li}_2 \left[\frac{1}{2}(1-z) \right] + \frac{1}{4} \ln^2 \left[\frac{1}{2}(1+z) \right] \\ - \frac{1}{2} \text{Li}_2 \left[\frac{1}{2}(1+z) \right] - \frac{1}{4} \ln^2 \left[\frac{1}{2}(1-z) \right]$$

with $\text{Li}_2(z)$ the dilogarithm function, defined by

$$\text{Li}_2(z) = - \int_0^z \xi^{-1} \ln(1-\xi) d\xi.$$

The quantity σ is the square root of a homogeneous six-degree polynomial in the α 's:

$$\begin{aligned} \sigma = & \left[\alpha_1^2 \alpha_{23}^2 (\alpha_1^2 - \alpha_2^2 - \alpha_3^2 - \alpha_{12}^2 + \alpha_{23}^2 - \alpha_{31}^2) \right. \\ & + \alpha_2^2 \alpha_{31}^2 (-\alpha_1^2 + \alpha_2^2 - \alpha_3^2 - \alpha_{12}^2 - \alpha_{23}^2 + \alpha_{31}^2) \\ & + \alpha_3^2 \alpha_{12}^2 (-\alpha_1^2 - \alpha_2^2 + \alpha_3^2 + \alpha_{12}^2 - \alpha_{23}^2 - \alpha_{31}^2) \\ & \left. + \alpha_1^2 \alpha_2^2 \alpha_3^2 + \alpha_1^2 \alpha_{12}^2 \alpha_{31}^2 + \alpha_2^2 \alpha_{23}^2 \alpha_{12}^2 + \alpha_3^2 \alpha_{31}^2 \alpha_{23}^2 \right]^{1/2}. \end{aligned}$$

The $\gamma_k^{(j)}$ ($k, j = 0, 1, 2, 3$) are homogeneous third-degree polynomials in the α 's, and $\beta_k^{(j)}$ is defined by

$$\beta_k^{(j)} = (\sigma - \gamma_k^{(j)}) / (\sigma + \gamma_k^{(j)}).$$

Caution 1:

The functions u and v are multiple valued, the formula (9) holds only if the branches of them are chosen correctly.

Function	Branch points
$\ln(z)$	$0, \infty$
$\text{Li}_2(z)$	$1, \infty$
$u(z)$	$0, 1, \infty$
$v(z)$	$1, -1, \infty$

If the branch cut for the logarithm is taken to run from 0 to ∞ along the negative real axis and the principal branch is chosen as

$$\ln(z) = \ln|z| + i \arg(z), \quad -\pi < \arg(z) < \pi,$$

then the branch cuts and the principal branches for $\text{Li}_2(z)$, $u(z)$ and $v(z)$ are determined.

The singularities of the expression (9) may occur at

$$\sigma = 0,$$

$$(\gamma_k^{(j)})^2 - \sigma^2 = 0 \quad \text{for} \quad k, j = 0, 1, 2, 3$$

and

$$\gamma_k^{(k)} + \gamma_k^{(0)} = 0 \quad \text{for} \quad k \neq 0.$$

There are 20 cases should be considered!

- True singularities **VS.** Cancelling singularities

Luckily, the situations for singularities may be reduced to

$$\sigma = 0,$$

$$\begin{cases} -\alpha_1 + \alpha_2 + \alpha_3 = 0, \\ \alpha_1 - \alpha_2 + \alpha_3 = 0, \\ \alpha_1 + \alpha_2 - \alpha_3 = 0, \end{cases}$$

$$\begin{cases} -\alpha_1 + \alpha_{12} + \alpha_{31} = 0, \\ \alpha_1 - \alpha_{12} + \alpha_{31} = 0, \\ \alpha_1 + \alpha_{12} - \alpha_{31} = 0, \end{cases}$$

and

$$\begin{cases} -\alpha_2 + \alpha_{23} + \alpha_{12} = 0, \\ \alpha_2 - \alpha_{23} + \alpha_{12} = 0, \\ \alpha_2 + \alpha_{23} - \alpha_{12} = 0, \end{cases}$$

$$\begin{cases} -\alpha_3 + \alpha_{31} + \alpha_{23} = 0, \\ \alpha_3 - \alpha_{31} + \alpha_{23} = 0, \\ \alpha_3 + \alpha_{31} - \alpha_{23} = 0. \end{cases}$$

Three obstacles appear when applying the closed-form formula (9) to generate the basic integral

$J(n_1, n_2, n_3, n_{12}, n_{23}, n_{31}; \alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31})$:

- determine the correct branches of u and v
- remove the canceling singularities analytically
- differentiate over the α 's, especially at the canceling singularities

The corresponding solving methods:

- Branch tracking
- Use some analytic techniques
- Develop a new recursive relations for the differentials

- Branch tracking

Consider the path $(1, 1, 1, \lambda, \lambda, \lambda)$ ($0 \leq \lambda \leq 1$) which runs from SRP $(1, 1, 1, 1, 1, 1)$ to ARP $(1, 1, 1, 0, 0, 0)$.

The correct expression for I around SRP is given by

$$I(1, 1, 1, \lambda, \lambda, \lambda) = 16\pi^3(1 - 3\lambda^2)^{-1/2} \\ \times [v(z_1(\lambda)) + 3v(z_2(\lambda)) + 3v(z_3(\lambda)) \\ + 9v(z_4(\lambda)) + 3u(z_5(\lambda))]$$

with the choices of the principal branches of u and v and the square root being positive imaginary.

Over the straight-line path from $\lambda = 1$ to $\lambda = 0$, there are two canceling singular points at $\lambda = \frac{1}{\sqrt{3}}$ and $\lambda = \frac{1}{2}$, so the following modified path will be used in the tracking

$$\lambda = x, \quad 1 \geq x \geq \frac{1}{\sqrt{3}} + \delta,$$

$$\lambda = \frac{1}{\sqrt{3}} + \delta e^{i\theta}, \quad -\pi \leq \theta \leq 0,$$

$$\lambda = x, \quad \frac{1}{\sqrt{3}} - \delta \geq x \geq \frac{1}{2} + \delta,$$

$$\lambda = \frac{1}{2} + \delta e^{i\theta}, \quad -\pi \leq \theta \leq 0,$$

$$\lambda = x, \quad \frac{1}{2} - \delta \geq x \geq 0,$$

where δ is a small positive real number.

After the complicated tracking, we find the correct expression of I at ARP

$$\begin{aligned}
 & I(1, 1, 1, \lambda, \lambda, \lambda) \\
 &= 16\pi^3(1 - 3\lambda^2)^{-1/2} \lim_{\epsilon \rightarrow 0^+} \left\{ -\pi^2 - i\pi \ln \left[\frac{1 + z_1(\lambda)}{1 - z_1(\lambda)} \right] \right. \\
 &\quad + v(z_1(\lambda)) + 3v(z_2(\lambda) - i\epsilon) + 3v(z_3(\lambda) + i\epsilon) \\
 &\quad \left. + 9v(z_4(\lambda)) + 3u(z_5(\lambda)) \right\}
 \end{aligned}$$

with also the choices of the principal branches of u and v , and the square root now, changing to be positive real.

Hence the value of $I(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31})$ in the neighborhood of ARP is given by

$$\begin{aligned}
 & I(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) \\
 &= \frac{16\pi^3}{\sigma} \lim_{\epsilon \rightarrow 0^+} \left\{ -\pi^2 - i\pi \ln \left[\frac{1 + \gamma_0^{(0)}/\sigma}{1 - \gamma_0^{(0)}/\sigma} \right] + \sum_{j=1}^3 u(\beta_0^{(0)} \beta_0^{(j)}) \right. \\
 & \quad + v(\gamma_0^{(0)}/\sigma) + \sum_{j=1}^3 v(\gamma_j^{(j)}/\sigma - i\epsilon) + \sum_{j=1}^3 v(\gamma_0^{(j)}/\sigma + i\epsilon) \\
 & \quad \left. + \sum_{j=0}^3 \sum_{\substack{k=1 \\ k \neq j}}^3 v(\gamma_k^{(j)}/\sigma) \right\}.
 \end{aligned}$$

We consider the general case, which tracks from the SRP $(1, 1, 1, 1, 1, 1)$ to an arbitrary one $(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31})$:

- **Step 1:** Consider the path $((\alpha_1 - 1)\lambda + 1, 1, 1, 1, 1, 1)$, where λ runs from $\lambda = 0$ to $\lambda = 1$ and this means we track from $(1, 1, 1, 1, 1, 1)$ to $(\alpha_1, 1, 1, 1, 1, 1)$.
- **Step 2:** Consider the path $(\alpha_1, (\alpha_2 - 1)\lambda + 1, 1, 1, 1, 1), \dots$

⋮

- **Step 6:** Consider the path $(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, (\alpha_{31} - 1)\lambda + 1)$, which means tracking from $(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, 1)$ to $(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31})$.

We should point out here that, using some mathematical techniques can simplify the whole tracking process to some extent. But it is still incredibly complicated.

For a point

$$\begin{aligned} & (\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31}) \\ & = (872, 78.08, 34.09, 12, 12.788, 10.34), \end{aligned}$$

the correct expression of I around this point is

$$I(\alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31})$$

$$= \frac{16\pi^3}{\sigma} \lim_{\epsilon \rightarrow 0_+} \left\{ \begin{aligned} &u(\beta_0^{(0)} \beta_0^{(1)} + i\epsilon) + u(\beta_0^{(0)} \beta_0^{(2)}) + u(\beta_0^{(0)} \beta_0^{(3)}) \\ &+ \sum_{j=0}^1 \sum_{k=0}^3 v(\gamma_k^{(j)} / \sigma) + v(\gamma_0^{(2)} / \sigma - i\epsilon) + v(\gamma_1^{(2)} / \sigma + i\epsilon) \\ &+ v(\gamma_2^{(2)} / \sigma + i\epsilon) + v(\gamma_3^{(2)} / \sigma + i\epsilon) + v(\gamma_0^{(3)} / \sigma + i\epsilon) \\ &+ v(\gamma_1^{(3)} / \sigma - i\epsilon) + v(\gamma_2^{(3)} / \sigma + i\epsilon) + v(\gamma_3^{(3)} / \sigma + i\epsilon) \end{aligned} \right.$$

$$\begin{aligned}
& -i\pi \ln \left[\frac{1 + \gamma_0^{(0)}/\sigma}{1 - \gamma_0^{(0)}/\sigma} \right] + i\pi \ln \left[\frac{1 + \gamma_1^{(0)}/\sigma}{1 - \gamma_1^{(0)}/\sigma} \right] + 2i\pi \ln \left[\frac{1 + \gamma_2^{(0)}/\sigma}{1 - \gamma_2^{(0)}/\sigma} \right] \\
& -i\pi \ln \left[\frac{1 + \gamma_0^{(1)}/\sigma}{1 - \gamma_0^{(1)}/\sigma} \right] + i\pi \ln \left[\frac{1 + \gamma_2^{(1)}/\sigma}{1 - \gamma_2^{(1)}/\sigma} \right] + i\pi \ln \left[\frac{1 + \gamma_3^{(1)}/\sigma}{1 - \gamma_3^{(1)}/\sigma} \right] \\
& + i\pi \ln \left[\frac{1 + \gamma_0^{(2)}/\sigma}{1 - \gamma_0^{(2)}/\sigma} \right] + i\pi \ln \left[\frac{1 + \gamma_2^{(2)}/\sigma}{1 - \gamma_2^{(2)}/\sigma} \right] - i\pi \ln \left[\frac{1 + \gamma_3^{(2)}/\sigma}{1 - \gamma_3^{(2)}/\sigma} \right] \\
& - i\pi \ln \left[\frac{1 + \gamma_0^{(3)}/\sigma}{1 - \gamma_0^{(3)}/\sigma} \right] - 2i\pi \ln \left[\frac{1 + \gamma_3^{(2)}/\sigma}{1 - \gamma_3^{(2)}/\sigma} \right] - \pi^2 \left. \vphantom{\frac{1 + \gamma_0^{(3)}/\sigma}{1 - \gamma_0^{(3)}/\sigma}} \right\},
\end{aligned}$$

and

$$\begin{aligned}
& I(872, 78.08, 34.09, 12, 12.788, 10.34) \\
& = 5.272897938711298 \times 10^{-4}.
\end{aligned}$$

λ	$I(1, 1, 1, \lambda, \lambda, \lambda)$
0.05	$3.63236729624140 \times 10^2$
0.15	$2.59682121353657 \times 10^2$
0.25	$1.93536770465417 \times 10^2$
0.45	$1.17605235633348 \times 10^2$
0.70	$7.08583266363307 \times 10$
0.95	$4.65491345483885 \times 10$
1.00	$4.31360835924724 \times 10$

The future(ongoing) work

- To calculate
 $J(n_1, n_2, n_3, n_{12}, n_{23}, n_{31}; \alpha_1, \alpha_2, \alpha_3, \alpha_{12}, \alpha_{23}, \alpha_{31})$
- The relativistic and QED corrections
- The five-body problems

*THANK YOU SO MUCH FOR
YOUR ATTENTION!*